

A UNIQUE OSCILLATION SOLUTION TO THE SOLAR NEUTRINO PROBLEM?

M. B. SMY

Department of Physics and Astronomy

4182 Frederick Reines Hall, University of California, Irvine, CA 92697-4575

A global two-neutrino oscillation fit combining Super-Kamiokande solar neutrino data with the solar neutrino rates measured by Homestake, Gallex/GNO, SAGE and SNO prefers a single allowed area, the Large Mixing Angle solution, at about 95% confidence level. The mass² difference between the two mass eigenstates is $\Delta m^2 \approx 3\text{--}25 \times 10^{-5} \text{eV}^2$, the mixing angle θ is $\tan^2 \theta \approx 0.21\text{--}0.67$.

1 Introduction

Neutrino flavor oscillations are currently the most favored way to understand the larger solar ^8B neutrino interaction rate measured by Super-Kamiokande¹ (SK) with respect to the ^8B rate reported by the SNO-collaboration². The SNO rate includes only e -type ^8B neutrinos while SK measures solar neutrinos via neutrino-electron elastic scattering which has a small sensitivity to other (active) flavors as well. The SK excess is then interpreted as a hint for appearance of other active flavors in the ^8B neutrino flux (born in the sun as purely e -type). Neutrino oscillations also explain the deficit of both rates as well as the deficits of the “Chlorine rate” (Homestake³) and the “Gallium rates” (Gallex/GNO⁴, SAGE⁵) with respect to the Standard Solar Model⁶ (SSM). Chlorine and Gallium rates include only e -type solar neutrinos.

The large neutrino flavor mixing between the second and third generation inferred from atmospheric neutrino data⁷ and the absence of an oscillation signal in the CHOOZ reactor neutrino experiment⁸ requires the neutrino flavor mixing between the first and third generation to be small. Solar neutrino oscillations can therefore be approximated by a two-neutrino description with the parameters θ (mixing angle) and Δm^2 (difference in mass²). The mixing angle ranges between 0 and $\pi/2$, since Δm^2 is defined to be positive. For Δm^2 between $\approx 10^{-8} \text{eV}^2$ and $\approx 10^{-3} \text{eV}^2$, the matter density in the sun and earth can strongly affect the oscillation probability. For $\theta < \pi/4$, resonant enhancement of the oscillations (MSW effect⁹) can occur in the sun. On the “dark side” of the parameter space ($\theta > \pi/4$), an anti-resonance can suppress the oscillations. Below $\approx 10^{-9} \text{eV}^2$ (quasi-vacuum/vacuum region), the oscillation probability is more affected by the oscillation phase than by matter effects.

2 Parameter Estimation for Two-Neutrino Oscillations

The simplest way to constrain the oscillation parameters is a χ^2 fit to all e -type neutrino rates, that is Gallex/GNO and SAGE (combined into a single “Gallium” number), Homestake and SNO. The dotted lines of figure 1a) show the allowed regions at 95% C.L. (which can be understood as an “overlay” of the allowed regions¹⁰ from the three experimental rates mentioned above and their SSM-based predictions). The area above $\Delta m^2 \approx 10^{-5} \text{eV}^2$ near maximal mixing is called the Large Mixing Angle (LMA) solution. The Small Mixing Angle Solution (SMA) is located between $\Delta m^2 \approx 10^{-6} \text{eV}^2$ and $\approx 10^{-5} \text{eV}^2$ at $\tan^2 \theta \approx 10^{-3}$. The LOW solution is the large Δm^2 part of the extensive region(s) between $\approx 10^{-7} \text{eV}^2$ and $\approx 10^{-9} \text{eV}^2$. The lower part of this region is called the quasi-vacuum (quasi-VAC) solution. The vacuum solutions (VAC) are below $\approx 10^{-10} \text{eV}^2$. All regions have similar χ^2 (LOW fits slightly worse).

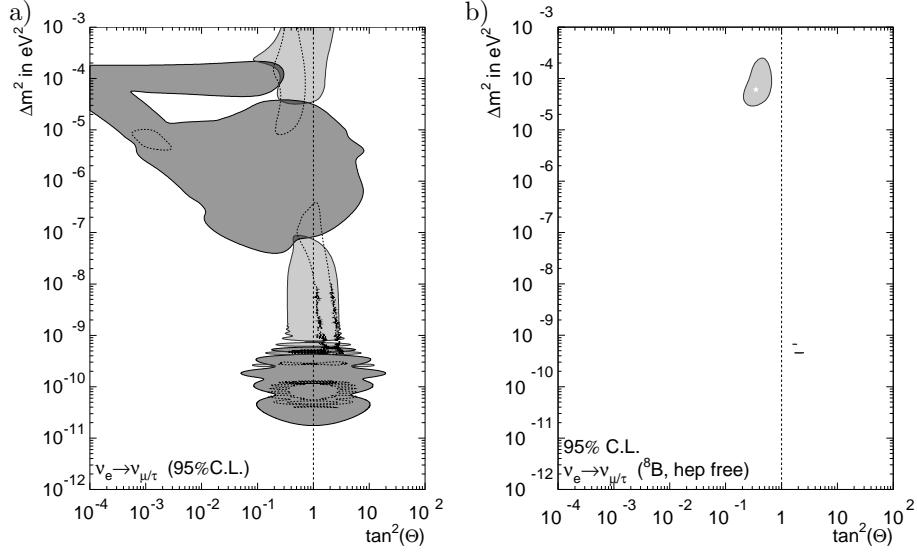


Figure 1. SK Zenith Spectrum/SNO (a) and global fit (b) contours at 95% C.L. The dark-grey area in a) is excluded by the shape of the SK Zenith Spectrum. Adding the SNO rate measurement results in the (light-grey) allowed areas. Both regions do not depend on neutrino flux predictions by the SSM. Overlaid are the allowed regions (inside dotted lines) based on the Gallex/GNO, SAGE, Homestake and SNO rates and the SSM. In b) the allowed areas from a global fit (adding the rates of Gallex/GNO, SAGE and Homestake to the SK Zenith Spectrum–SNO fit) are predominantly LMA solutions.

To break the degeneracy in χ^2 of these regions, more information is needed than provided by the suppression of the e -type rates. Indeed, neutrino oscillations not only cause a ‘disappearance’ of the e -type fluxes but also produce other neutrino flavors (appearance). Furthermore, they can induce distortions of the neutrino spectrum and time variations of the solar neutrino fluxes. In the MSW region, the time variations arise from matter effects inside the earth (daily variations). In the vacuum regions, the time variations are the consequence of the change of the oscillation phase due to the yearly variation of the oscillation baseline (distance between sun and earth).

SK can search for spectral distortions through analysis of the spectrum of the recoiling electron as well as daily or yearly time variations. The analysis of the shape of the SK “zenith angle spectrum”^{10,11} combines spectrum and daily variation analyses. Neither spectral distortion nor daily variation was found. The centerpiece of this work is the preliminary SK zenith angle spectrum based on the entire SK data set (1496 days). The numerical results are listed in tables 1 and 2. Figure 1a) shows the excluded regions (dark-grey) from the SK zenith angle spectrum (for an explanation of the fit method see Appendix, eq. 2). The SMA and VAC regions are excluded since they require a distorted ^8B neutrino spectrum¹⁰. The lower part of the LMA and the upper part of the LOW solution predict daily variations¹⁰ and are therefore disfavored. The excluded areas are independent of the SSM neutrino fluxes. The SK zenith angle spectrum shape is best described by quasi-VAC solutions.

If combined with the SNO rate² of $0.346^{+0.029}_{-0.028} \times \text{SSM}$ (Appendix, eq. 3), the SK rate provides a probe for the appearance of other neutrino flavors. The SK rate of $0.465^{+0.015}_{-0.013} \times \text{SSM}$ (see table 1) exceeds the e -type rate inferred from the SNO measurement by more than 3σ . If this is interpreted as appearance, the other flavors contribute about 25% to the SK rate and 70% to the ^8B flux (the SK cross section for the other flavors is six to seven times smaller than for e -type neutrinos). The two light-grey allowed regions of figure 1a) are based on the combined fit to the SK zenith angle spectrum and the SNO rate¹². One region contains the upper part of the LMA solution, the other region contains the lower part of the LOW solution and the quasi-VAC solutions (best-fit). The allowed areas are still independent of the SSM neutrino fluxes.

The Appendix (eq. 4) explains the method to add the Gallium^{4,5} ($74.8^{+5.1}_{-5.0} \text{SNU}$ or $0.584^{+0.040}_{-0.039} \times \text{SSM}$) and Chlorine rates³ ($2.56 \pm 0.23 \text{SNU}$ or $0.337 \pm 0.030 \times \text{SSM}$) to the fit. The allowed areas of this global fit (shown in figure 1b) look quite different when compared with the e -type rate fit: Only the upper part of the LMA survives (and two tiny quasi-VAC solutions). The disappearance of the LOW solution is due to predicted daily variations. The SK rate requires a larger ^8B flux than the Chlorine rate allows, so almost all

Table 1. SK Rate and Uncertainty for eight energy bins. The rates, statistical and systematic uncertainties (of the spectrum shape) in the third column are presented in units of SSM expectation. These systematic uncertainties are assumed to be uncorrelated in energy. The uncertainties in the fourth (uncertainty of the ^8B neutrino spectrum), fifth (uncertainty of the energy scale of 0.64%) and sixth column (uncertainty of the energy resolution of 2.5%) are fully correlated in energy (but uncorrelated with each other). The combined uncertainty (last row) is based on the sum of all bins of the zenith angle spectrum. The combined rate has an additional systematic uncertainty of $\sigma_{SK} = {}^{+2.9}_{-2.6}\%$ (excluding ^8B shape, energy scale and resolution), which was added to the uncorrelated systematic uncertainties of the spectrum shape.

Bin	Range [MeV]	Rate \pm stat \pm syst [SSM]	^8B Spectrum	E-Scale	E-Resol.
1	5.0-5.5	$0.4671\pm 0.0404 {}^{+0.0165}_{-0.0138}$	$+0.04\%$ -0.02%	$+0.09\%$ -0.01%	$+0.23\%$ -0.21%
2	5.5-6.5	$0.4580\pm 0.0141 {}^{+0.0066}_{-0.0065}$	$+0.13\%$ -0.09%	$+0.20\%$ -0.16%	$+0.19\%$ -0.17%
3	6.5-8.0	$0.4729\pm 0.0084\pm 0.0065$	$+0.41\%$ -0.38%	$+0.63\%$ -0.62%	$+0.17\%$ -0.16%
4	8.0-9.5	$0.4599\pm 0.0093\pm 0.0063$	$+0.89\%$ -0.85%	$+1.3\%$ -1.3%	$+0.12\%$ -0.12%
5	9.5-11.5	$0.4627\pm 0.0103\pm 0.0063$	$+1.7\%$ -1.6%	$+2.5\%$ -2.4%	$+0.16\%$ -0.18%
6	11.5-13.5	$0.4621\pm 0.0168\pm 0.0063$	$+3.1\%$ -2.7%	$+4.4\%$ -4.1%	$+1.1\%$ -1.1%
7	13.5-16.0	$0.5666\pm 0.0390\pm 0.0078$	$+5.1\%$ -4.2%	$+7.0\%$ -6.4%	$+3.2\%$ -3.2%
8	16.0-20.0	$0.5554\pm 0.1458\pm 0.0076$	$+7.7\%$ -5.6%	$+10.6\%$ -9.6%	$+8.4\%$ -7.9%
Comb.	5.0-20.0	$0.4653\pm 0.0047 {}^{+0.0138}_{-0.0122}$	$+1.15\%$ -1.04%	$+1.66\%$ -1.58%	$+0.33\%$ -0.34%

Table 2. Subdivision of bins 2-7 according to the solar zenith angle θ_z . The range of $\cos\theta_z$ is given for each bin: $\cos\theta_z < 0$ is ‘Day’ and $\cos\theta_z > 0$ is ‘Night’ (‘Mantle’ and ‘Core’). The rates are given in units of $0.001\times\text{SSM}$. Only statistical uncertainties are quoted. All systematic uncertainties (see table 1) are assumed to be fully correlated in zenith angle.

Bin	Day	Mantle					Core
	-1-0	0.00-0.16	0.16-0.33	0.33-0.50	0.50-0.67	0.67-0.84	0.84-1
2	453 ± 20	442 ± 53	379 ± 49	472 ± 45	522 ± 45	503 ± 49	426 ± 52
3	474 ± 12	530 ± 34	506 ± 30	438 ± 26	478 ± 26	451 ± 28	439 ± 31
4	448 ± 13	463 ± 36	470 ± 33	462 ± 29	509 ± 29	461 ± 32	451 ± 35
5	453 ± 15	449 ± 40	502 ± 38	451 ± 32	473 ± 32	477 ± 35	483 ± 40
6	477 ± 25	509 ± 67	351 ± 55	391 ± 49	498 ± 53	434 ± 56	521 ± 64
7	511 ± 54	570 ± 150	831 ± 167	694 ± 131	665 ± 127	441 ± 118	469 ± 131

the quasi-VAC regions disappear. The fit does not depend on the ^8B and *hep* neutrino flux predictions by the SSM, which suffer from the largest uncertainties. However, it depends on the other SSM neutrino fluxes, in particular the ^7Be flux (10% uncertainty) and the neutrino fluxes of the CNO cycle ($\approx 20\%$ uncertainty). Those fluxes contribute⁶ about 15% (^7Be) and 6% (CNO) to the Chlorine rate and 27% (^7Be) and 7% (CNO) to the Gallium rate in the SSM.

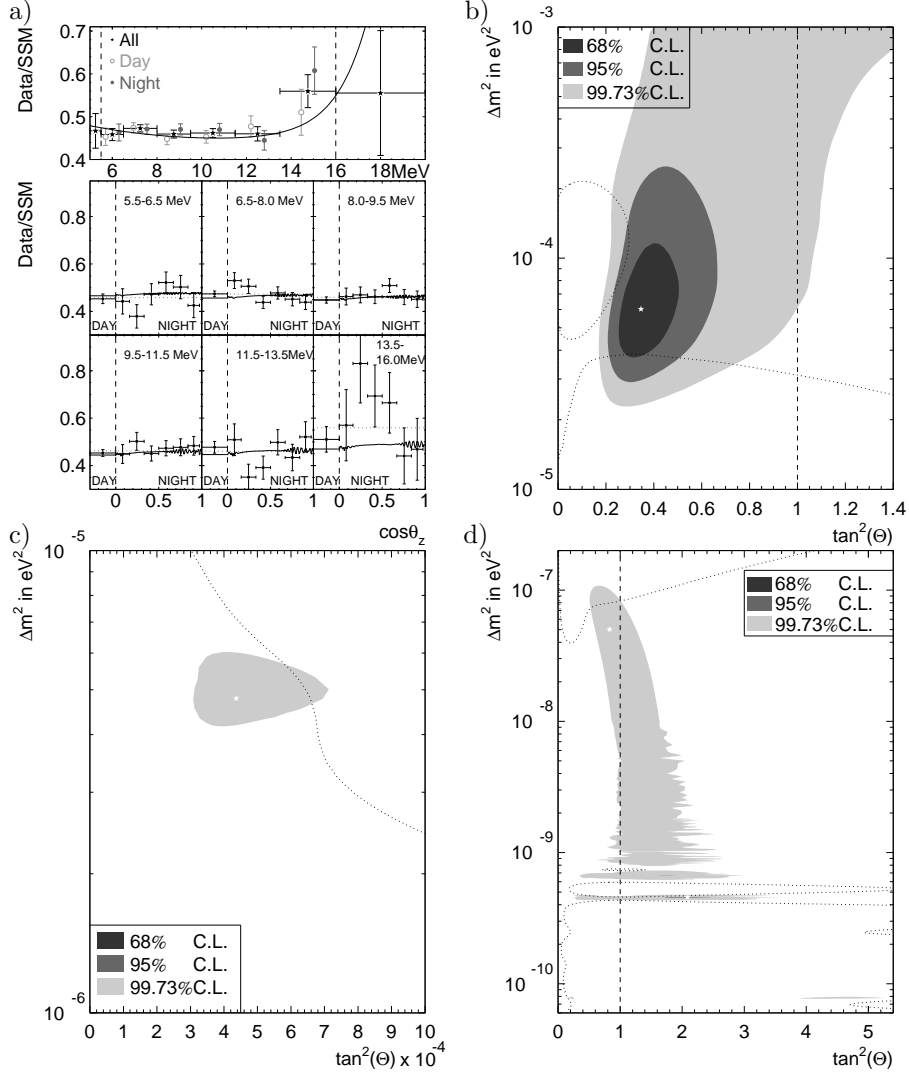
Table 3. Parameters for the best-fit points. The probabilities given in the fourth row are based on the difference of χ^2 (with respect to the minimum). The five rows below show the five independent parts of the fit: the $\Delta\chi^2$ from a fit to the shape of the SK zenith spectrum (Appendix, eq. 2) and four interaction rates (deviation probabilities are given in units of Gaussian standard deviation σ). The last four rows show the values of the minimized fit parameters. The ^8B and the *hep* fluxes are free, the ^8B neutrino spectrum shift as well as the SK energy scale and resolution shifts are constrained within the systematic uncertainties.

Solution	Large Mixing Angle (LMA)	Quasi-Vacuum (Quasi-VAC)	Low Δm^2 (LOW)	Small Mixing Angle (SMA)
Δm^2	6.0×10^{-5}	4.57×10^{-10}	5.0×10^{-8}	4.8×10^{-6}
$\tan^2 \theta$	0.35	2.1	0.83	0.00044
χ^2 (45 dof; p_{χ^2} [%])	43.4 (53.9)	48.5 (33.4)	51.2 (24.3)	54.2 (16.2)
$\Delta\chi^2$ (2 dof; $p_{\Delta\chi^2}$ [%])	0.0 (100.0)	5.1 (7.9)	7.8 (2.0)	10.8 (0.5)
$\Delta\chi^2_{SK}$ ($p_{\Delta\chi^2}$ [%])	3.4 (1.3 σ)	3.1 (1.3 σ)	3.9 (1.5 σ)	5.0 (1.7 σ)
Ga Rate [SNU]	73.2 (−0.3 σ)	69.6 (−1.0 σ)	68.2 (−1.3 σ)	75.1 (+0.1 σ)
Cl Rate [SNU]	2.97 (+1.8 σ)	3.18 (+2.8 σ)	3.13 (+2.5 σ)	2.67 (+0.5 σ)
SK Rate [%SSM]	46.4 (−0.1 σ)	44.7 (−1.4 σ)	44.9 (−1.2 σ)	44.1 (−1.9 σ)
SNO Rate [%SSM]	32.8 (−0.7 σ)	37.1 (+0.8 σ)	38.5 (+1.3 σ)	43.8 (+3.1 σ)
ϕ_{s_B} [$10^6/(\text{cm}^2\text{s})$]	5.62 (+0.6 σ)	3.71 (−1.7 σ)	4.04 (−1.2 σ)	2.71 (−2.9 σ)
ϕ_{hep} [$10^3/(\text{cm}^2\text{s})$]	40	0	21	8
^8B Spectrum Shape	−0.3 σ	−0.7 σ	−0.1 σ	+0.1 σ
SK E-scale/resol.	−0.4 σ /−0.1 σ	−1.0 σ /0.0 σ	−0.1 σ /−0.2 σ	+0.1 σ /−0.4 σ

3 Results of the global fit

Table 3 compares the four smallest local minima of the χ^2 describing the fit. The best-fit is located in the upper LMA area. The ^8B flux resulting from this fit is somewhat higher than expected by the SSM ($5.05^{+1.01}_{-0.81} \times 10^6/(\text{cm}^2\text{s})$) but well within the uncertainty. The *hep* flux is considerably higher than expected by the SSM ($9.3 \times 10^3/(\text{cm}^2\text{s})$); however, the uncertainty of this prediction is thought to be very large. The fit agrees with the SK zenith angle spectrum moderately well. Figure 2a) shows the SK zenith angle spectrum and this best fit. Figure 2b) gives a magnified view of the LMA region (in a linear scale of $\tan^2 \theta$). The 1 σ (black) and 3 σ (light grey) contours are also shown. The fit accommodates well the Gallium, SK and SNO rates, however the predicted Chlorine rate is about 2 σ too high. This worsens the otherwise very good best-fit χ^2 . As a consequence, quasi-VAC solutions, which don't fit well the interaction rates in general, appear at 95% C.L.

The quasi-VAC solution describes the SK zenith angle spectrum slightly better than the LMA. The fit is helped by a statistical up-fluctuation of the last two bins (see figure 3). Together with a down-shift of energy scale and ^8B spectrum, the required oscillation minimum can be ‘generated’ in the data.



The rates, however, do not fit well. The fit struggles to accommodate both the Chlorine and SK rate. In this region Chlorine and SK disagree by 3σ about the ^8B flux (which in the un-oscillated SSM contributes about 76% of the Chlorine rate). The resulting best-fit ^8B flux falls 1.7σ short of the SSM prediction. Even though the quasi-VAC solution has an overall ‘C.L. threshold’ of 8%, it is considerably disfavored, when thus checked in detail. The surviving (at 98% C.L.) LOW solution fits the rates about as poorly as the quasi-VAC solution. The SK zenith angle spectrum fits somewhat worse than either LMA or quasi-VAC. The lack of zenith angle variation in the SK data reduces the Δm^2 (usually around 10^{-7}eV^2) and worsens the LOW best-fit which is already under pressure from the rates. Figure 2d) shows an enlarged view of the LOW and quasi-VAC region.

A “smaller-than-small mixing angle” solution appears just left of the SMA at about the 3σ level (see figure 2c). The SMA region is defined by the crossing of the Gallium and the Chlorine allowed area (for a given ^8B flux). It therefore fits those two rates very well. The SK zenith angle spectrum, however, fits the worst of all solutions: The SK spectral data lack the predicted distortion. The SNO rate is more than 3σ above the measurement, the SK rate is too low (2σ). The ^8B flux required is very low (3σ).

The ^8B flux comparisons above refer to the SSM value. A recent precision measurement of the cross section of the $^7\text{Be}(p,\gamma)^8\text{B}$ fusion reaction by Junghans et al.¹³ implies (see ¹⁴) a flux of $5.93^{+0.83}_{-0.89} \times 10^6/(\text{cm}^2\text{s})$. If the fit is confined to this flux, then all solutions other than LMA are further disfavored.

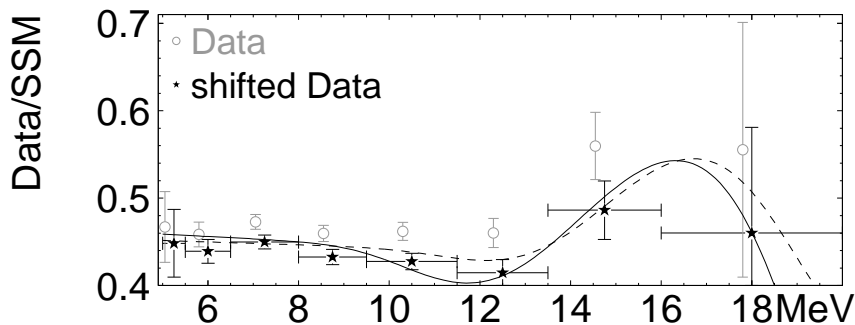


Figure 3. SK Spectral distortion for two quasi-VAC solutions (solid line: best fit at $4.57 \cdot 10^{-10}\text{eV}^2$, dashed line: $6.68 \cdot 10^{-10}\text{eV}^2$) and data (open circles). Although the data do not support spectral distortions, a shift (stars) in energy scale (-1.0σ) and ^8B shape (-0.7σ) and systematic uncertainty of the combined rate (-1.4σ) can accommodate the prediction.

4 Conclusion

A global fit to all solar neutrino data was performed using a two-neutrino oscillation model. At 95% C.L. only the upper LMA solution and two tiny quasi-VAC regions are still allowed. The quasi-VAC regions are disfavored at more than 92% C.L. caused by a disagreement about the ^8B flux between SK and the Chlorine rate of 3σ ; the resulting quasi-VAC ^8B flux fit is 1.7σ smaller than the Standard Solar Model. The LMA solutions above $\Delta m^2 = 3 \cdot 10^{-5} \text{eV}^2$ are therefore the only viable solution at 95% C.L.

Acknowledgments

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Appendix: Fitting Method

To estimate and limit oscillation parameters, a χ^2 describing the shape of the SK zenith spectrum is extended to take into account the measured neutrino interaction rates by various experiments. Using oscillation probabilities (obtained as in ¹¹) and SSM neutrino fluxes, the expected interaction rates $B_{i,z}^{\text{osc}}$ (due to ^8B neutrinos) and $H_{i,z}^{\text{osc}}$ (due to *hep* neutrinos) in energy bin i and zenith-angle bin z are calculated. These rates as well as the SK measurements $D_{i,z}$ are then normalized by the SSM expectations without oscillations:

$$b_{i,z} = \frac{B_{i,z}^{\text{osc}}}{B_{i,z}^{\text{SSM}} + H_{i,z}^{\text{SSM}}}, \quad h_{i,z} = \frac{H_{i,z}^{\text{osc}}}{B_{i,z}^{\text{SSM}} + H_{i,z}^{\text{SSM}}}, \quad d_{i,z} = \frac{D_{i,z}}{B_{i,z}^{\text{SSM}} + H_{i,z}^{\text{SSM}}}. \quad (1)$$

The zenith components of the flux difference vector $\vec{\Delta}_i$

$$\Delta_{i,z}(\beta, \eta) = (\beta \cdot b_{i,z} + \eta \cdot h_{i,z}) \times f(E_i, \delta_B, \delta_S, \delta_R) - d_{i,z}$$

allow for arbitrary neutrino fluxes (through the free parameters β and η). The combined rate prediction is modified by the energy-shape factors

$$f(E_i, \delta_B, \delta_S, \delta_R) = f_B(E_i, \delta_B) \times f_S(E_i, \delta_S) \times f_R(E_i, \delta_R)$$

with δ_B describing the ^8B neutrino spectrum shape uncertainty, δ_S describing the uncertainty of the SK energy scale (0.64%) and δ_R describing the

uncertainty of the SK energy resolution (2.5%) (The shapes f_B, f_S, f_R are given in table 1). All three uncertainties affect the bins of the SK zenith angle spectrum in a correlated way. The 7×7 matrices V_i describe statistical and energy-uncorrelated uncertainties; the latter are assumed to be fully correlated in zenith angle. For any given parameters δ_k , the χ^2

$$\chi_0^2 = \sum_{i=1}^8 \vec{\Delta}_i \cdot V_i^{-1} \cdot \vec{\Delta}_i = \chi_{0,m}^2 + \left(\vec{\phi} - \vec{\phi}_{0,m} \right) C_0 \left(\vec{\phi} - \vec{\phi}_{0,m} \right)$$

$$\text{with } C_0 = \sum_{i=1}^8 \begin{pmatrix} \vec{b}_i \cdot V_i^{-1} \cdot \vec{b}_i & \vec{h}_i \cdot V_i^{-1} \cdot \vec{b}_i \\ \vec{h}_i \cdot V_i^{-1} \cdot \vec{b}_i & \vec{h}_i \cdot V_i^{-1} \cdot \vec{h}_i \end{pmatrix} \quad \text{and} \quad \vec{\phi} = \begin{pmatrix} \beta \\ \eta \end{pmatrix}$$

can be written as a quadratic form of $\vec{\phi}$ and the curvature matrix C_0 . The minimum is

$$\chi_{0,m}^2 = \sum_{i=1}^8 \vec{d}_i \cdot V_i^{-1} \cdot \vec{d}_i - C_{0,m} \quad \text{with} \quad C_{0,m} = \vec{\phi}_{0,m} C_0 \vec{\phi}_{0,m}.$$

If the minimum flux vector is scaled by α ($\vec{\phi} = \alpha \times \vec{\phi}_{0,m}$) then χ_0^2 constrains α to be $\alpha = 1 \pm \sqrt{1/C_{0,m}}$. To take into account the systematic uncertainty of the SK combined rate $\sigma_{SK} = {}^{+2.9}_{-2.6} \%$ (which is completely correlated in zenith-angle and energy), χ_0^2 is modified to

$$\chi_1^2 = \chi_{0,m}^2 + \left(\vec{\phi} - \vec{\phi}_{0,m} \right) C_1 \left(\vec{\phi} - \vec{\phi}_{0,m} \right) \quad \text{with} \quad C_1 = \frac{1/\sigma_{SK}^2}{C_{0,m} + 1/\sigma_{SK}^2} \times C_0.$$

χ_1^2 constrains α to be $\alpha = 1 \pm \sqrt{1/C_{0,m} + \sigma_{SK}^2}$, that is, the minimum is unchanged, but the allowed parameter range for β and η is larger. The total χ^2 for the SK zenith spectrum shape is then

$$\chi_{SK}^2 = \text{Min} \left(\chi_1^2(\beta, \eta, \delta_B, \delta_S, \delta_R) + \left(\frac{\delta_B}{\sigma_B} \right)^2 + \left(\frac{\delta_S}{\sigma_S} \right)^2 + \left(\frac{\delta_R}{\sigma_R} \right)^2 \right) \quad (2)$$

where all δ_k as well as β, η are minimized.

A combined fit with the rate measured by the SNO collaboration² can be done by defining b_{SNO}, h_{SNO} and $\Delta_{SNO}(\beta, \eta)$ in a similar way as in (1) and form

$$\chi_{SK-SNO}^2 = \text{Min} \left(\chi_1^2 + \left(\frac{\delta_B}{\sigma_B} \right)^2 + \left(\frac{\delta_S}{\sigma_S} \right)^2 + \left(\frac{\delta_R}{\sigma_R} \right)^2 + \left(\frac{\Delta_{SNO}(\beta, \eta)}{\sigma_{SNO}} \right)^2 \right) \quad (3)$$

To add the radio-chemical rate measurements of Homestake³, Gallex/GNO⁴, and SAGE⁵ (all “Gallium” rates are combined into a single rate), the χ^2

$$\chi_{\text{RC}}^2(\beta, \eta)$$

takes into account the correlations between Gallium and Chlorine measurements. The ^8B and *hep* fluxes are constrained by the minimization of $\chi_{\text{SK-SNO}}^2$, not by the Standard Solar Model. The total χ^2 is then a simple addition

$$\chi^2 = \chi_{\text{SK-SNO}}^2 + \text{Min} \left(\chi_{\text{RC}}^2(\alpha\beta_{\text{min}}, \alpha\eta_{\text{min}}) + \left(\frac{\alpha - 1}{\sigma_\alpha} \right)^2 \right) \quad (4)$$

where $1/\sigma_\alpha^2 = C_{\text{SK-SNO,m}} = \vec{\phi}_{\text{SK-SNO,m}} \cdot C_{\text{SK-SNO}} \cdot \vec{\phi}_{\text{SK-SNO,m}}$. Numerically, σ_α is found to be about 2.7% (The accuracy of SK and SNO combined)

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